"Result of search for solar hadronic axions emitted in M1-transition of ⁸³Kr"

Akhmatov⁴ Z.A., Yu.M. Gavrilyuk¹, A.M. Gangapshev¹, A.V. Derbin³, I.S. Drachnev³, V.V. Kazalov¹, Khokonov A.Kh-A.⁴, V.V. Kuzminov¹, V.N. Muratova, S.I. Panasenko², S.S. Ratkevich², D.A. Tekueva¹, S.P. Yakimenko¹, E.V. Unzhakov³

¹ – Institute for Nuclear Research RAS, Moscow, Russia
 ² – Kharkov National University, Ukraine
 ³ – St. Petersburg Nuclear Physics Inst., NRCKI ,Gatchina, Russia
 ⁴-Kabardino-Balkarian State University, Nalchik, Russia
 Work is supported by RFBR: 14-02-00258 and 16-29-13011



ICSSNP-2017 "Physics of Fundamental Interactions"

"Axions are among the most fascinating particles on the long list of those proposed but not yet observed or ruled out. Their existence would provide an elegant resolution of the strong CP problem. Even more exciting is the possibility that the missing mass needed to close the universe is composed of axions, and that axions are «cold dark matter» which seems to be necessary for galaxy formation. ..."

Mark Srednicki, "Axion couplings to matter (I). CP-conserving parts", Nucl. Phys. B260 (1985) 689-700.

"...the composite axion is a particular example of a "hadronic" axion, resulting from a theory where only exotic fermions carry $U(1)_{PQ}$ charges. Hadronic axions don't couple to leptons, which are neutral under $SU(3)xU(1)_{PQ}$. Nor do they couple to heavy quarks, which are integrated out of the theory above 1GeV, where QCD gets strong. Hadronic axions will still couple to nucleons as well as to photons. ..."

David B. Kaplan, "Opening the axion window", Nucl. Phys. B260 (1985) 215-226.

"The most attractive solution of the strong CP problem is to introduce the Peccei-Quinn global symmetry which is spontaneously broken at energy scale f_a . The original axion model assumed that f_a is equal to the electroweak scale. Although it has been experimentally excluded, variant "invisible" axion models are still viable in which f_a is assumed to be very large. ... Such models are referred to as hadronic and Dine-Fischler-Srednicki-Zhitnitskii axions." Shigetaka Moriyama, "Proposal to search for a monochromatic component of solar axions using E_{a} .

⁵⁷Fe", Phys. Rev. Lett. v.75 №8 (1995) 3222-3225.

The axions canbe produced when thermally excited nuclei (or excited due to nuclear reactions) in the Sun relaxes by magnetic transition to its ground state and could be detected via resonant excitation of the same nuclide in a laboratory.



The monochromatic axions emitted by ⁷Li, ⁵⁷Fe and ⁸³Kr nuclei can excite the same nuclide in a laboratory, because the axions are Doppler broadened due to thermal motion of the axion emitter in the Sun, and thus some axions have needed energy to excite the nuclide.

Solar axions spectra vs g_{Ay} , g_{Ae} and g_{AN}



The main sources of solar axions: 1. Reactions of main solar chain. The most intensive fluxes are expected from M1transitions in ⁷Li and ³He nuclei(g_{AN}): ⁷Be+e⁻ \rightarrow Li*+ γ ; ⁷Li* \rightarrow ⁷Li+A(478кэB) $p + d \rightarrow {}^{3}He + A (5.5 M \rightarrow B).$ 2. Magnetic type transitions in nuclei whose low-lying levels are excited due to high temperature in the Sun(57 Fe, 83 Kr)(g_{AN}) 3. Primakoff conversion of photons in the electric field of solar $plasma(g_{Ay})$. 4. Bremsstrahlung: $e+Z(e) \rightarrow Z+A.(g_{Ae})$ 5. Compton process : $\gamma + e \rightarrow e + A.(g_{Ae})$ 6. axio-recombination: $e + I \rightarrow I^- + A$ and axio-deexcitation: $I^* \rightarrow I + A$. PRD 83 023505 (2011) CAST 1302.6283, 1310.0823

Flux of solar ⁸³Kr axions

The total axion flux Φ_A depends on the level energy $E_{\gamma} = 9.4$ keV, temperature T, nuclear level lifetime $\tau_{\gamma} = 3.6$ µs, the abundance of the ⁸³Kr isotope on the Sun N, and the branching ratio of axions to photons emission ω_A/ω_{γ} :

$$\Phi_{A} = \int N(r) \frac{2 \exp(-E_{\gamma} / kT(r))}{1 + 2 \exp(-E_{\gamma} / kT(r))} \frac{\omega_{A}}{\tau_{\gamma} \omega_{\gamma}} dr$$
$$\Phi_{A}(E_{M1}) = 5.97 \times 10^{23} \left(\frac{\omega_{A}}{\omega_{\gamma}}\right) \text{cm}^{-2} \text{s}^{-1} \text{keV}^{-1}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

$$m_{a} = \frac{\sqrt{z}}{1+z} \frac{f_{\pi}m_{\pi}}{f_{a}} = 1 \text{eV} \frac{\sqrt{z}}{1+z} \frac{1.3 \cdot 10^{7}}{f_{a}/\text{GeV}}$$

Resonant absorption by ⁸³Kr nucleus

$$\sigma\left(E_{A}\right) = 2\sqrt{\left(\pi\right)}\sigma_{0\gamma}\exp\left[\frac{-4\left(E_{A}-E_{M1}\right)}{\Gamma^{2}}\right]\frac{\omega_{A}}{\omega_{\gamma}}$$
$$\frac{\omega_{A}}{\omega_{\gamma}} = \frac{1}{2\pi\alpha}\frac{1}{1+\delta^{2}}\left[\frac{g_{AN}^{0}\beta+g_{AN}^{3}\beta}{\left(\mu_{0}-0.5\right)\beta+\mu_{3}-\eta}\right]^{2}\left(\frac{p_{A}}{p_{\gamma}}\right)^{3}$$

where δ – quenching factor for M1 transition, $\alpha \approx$ 1/137, $\mu_0 \approx 0.88$ and $\mu_3 \approx$ 4.71– izoscalar and izovector nuclear magnetic moments, β =-1 and η =0.5 – nuclear matrix constants.

$$g_{AN}^{0} = -7.8 \cdot 10^{-8} \left(\frac{6.2 \cdot 10^{6}}{f_{a}/GeV} \right) \left(\frac{3F - D + 2S}{3} \right) \qquad g_{AN}^{3} = -7.8 \cdot 10^{-8} \left(\frac{6.2 \cdot 10^{6}}{f_{a}/GeV} \right) \left((D + F) \frac{1 - z}{1 + z} \right)$$

$$R[g^{-1}day^{-1}] = 4.23 \times 10^{21} \left(\frac{\omega_A}{\omega_\gamma}\right)^2 =$$

=8.53×10²¹ $\left(g^3_{AN} - g^0_{AN}\right)^4 \left(\frac{p_A}{p_\gamma}\right)^6 =$
=2.41×10⁻¹⁰ $m^4_A \left(\frac{p_A}{p_\gamma}\right)^6$

Experimental Setup

Detector	Proportional counter
Shield	23cm Pb, 8cm PE, 20cm Cu
Working gas	99.9% Kr-93
Diameter (case)	134 mm
Anode (gold coated tungsten)	10 μm
Fidutial length	595 mm
Gas pressure	1.8 at





PA – *Preamplifier* A1 – Amplifier DO – FADC (Digitizer)

Phase II – 613 days spectrum of 99% ⁸³Kr-detector



Spectra of Kr-chamber measured during 613 days of livetime. The peak of 13.5keV from K-capture of ⁸¹Kr is in 1.2×10^3 times less intensive than PhaseI results. Since the proportional chamber is prepared from cupper and Cu X-rays are clear visible.

Results:



 $|g_{AN}^{3} - g_{AN}^{0}| \le 8.4 \times 10^{-7},$ m_A $\le 65 \text{ eV}$ at 95% C.L.

Nuclide	Abundance in heliosphere, $\log_{10} s$ ($\log_{10} s(H)=12.00$)	1-st excitation level, keV (transition type)
Fe-57	5,78	14,41 (M1+E2)
Kr-83	2,34	9,40 (M1+E2)
Li-7	0,42	477,6 (M1+E2)

Overview of other experimental results

Article (nuclide)		Результат
1	M. Krcmar et al., arXive:nucl-ex/9801005v2 (Fe-57)	≤745 eV
2	K. Jakovčić et al., arXive:nucl-ex/0402016v1 (Kr-83)	≤5,5 keV
3	A.V. Derbin et al., Eur. Phys. J. C (2009) 62:755-760 (Fe-57)	≤159 eV
4	F.A. Danevich et al., arXive:nucl-ex/0811383v2 (Fe-57)	≤1,6 keV
5	P. Belli et al., Nucl. Phys. A (2008) 806:388-397 (Li-7)	≤13,9 keV
6	Yu.M. Gavrilyuk, et al., JETP Letters (2015,) (Kr-83)	≤99,5 eV
7	This work (Kr-81)	≤65 eV
	W.C. Haxton and K.Y. Lee, Phys. Rev. Lett. (1991)	≈3÷20eV

New detector for Solar hadronic axions

Detector: semiconductor detector based on Pyrite

```
Hall mobility: 0.5÷3 cm<sup>2</sup>/V*s
```

```
Electron mobility: 10÷50 cm<sup>2</sup>/V*s
```

Isotope abundance: ⁵⁴Fe-5,845%, ⁵⁶Fe-91,754%,

```
<sup>57</sup>Fe-2,119%, <sup>58</sup>Fe-0,282%

R[g^{-1}day^{-1}] = 8.48 \times 10^{-7} m_A^4 \left(\frac{p_A}{p_Y}\right)^6

\frac{R_{Fe-57}}{R_{Kr-83}} = 3.51 \times 10^3
```



goal:
$$S(m_A) \simeq 1 eV$$

Thank you for attention!

Pyrite



The mineral pyrite, or iron pyrite, also known as fool's gold, is an iron sulfide with the chemical formula FeS_2 . Pyrite has been proposed as an abundant, inexpensive material in low-cost photovoltaic solar panels. Synthetic iron sulfide was used with copper sulfide to create the photovoltaic material.



The band gap in pyrite is about 0.95 eV and the dominant charge carriers can be either electrons or holes. Sometimes, both n-type and p-type semiconducting regions can be found within single naturally occurring crystals

Monocrystall has a cubic form.

Resistivity (natiral crystalls): 10⁻⁵÷10⁰Om*m

Рабочие параметры сигналов



Распределение событий



Спектр после отбора событий (Живое время измерений 6585 часов, прамаетры отбора событий: Фронт<51кан, λ>109 и λ<280)

